

Nonlocal gravity

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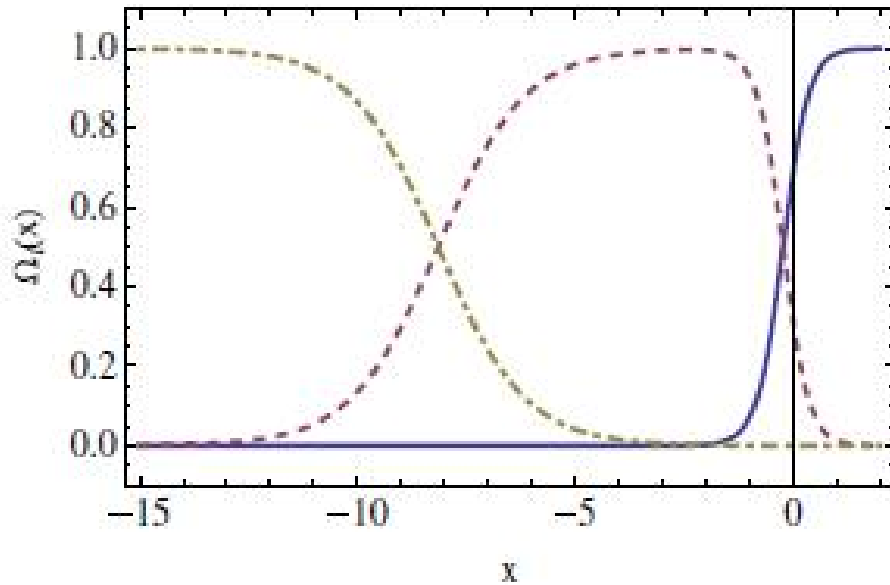
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Refs: PRD 98, 084040 (2018), PRD 99, 064044 (2019)

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- Motivation to modify gravity: The late-time acceleration
- General result to explain the acceleration



Maggiore et al, PRD, 2014

- Coincidence problem

The following may be a better theory:

- The late-time acceleration
- Newtonian potential $\Psi = \Phi$ exactly
- Except $\kappa = 8\pi G/c^2$, only dimensionless parameters appear in the Lagrangian
- Pure geometric theory

For dimensionless parameters:

- Dimension analysis: $[R]=\text{length}^{-2}$
- If we use R^2 , then we need $\text{length}^2 * R^2$
- The coefficient length^2 can be a parameter (c/H_0 etc), or a math operator (\Box^{-1} , an integral operator)
- $\Box^{-1} \rightarrow$ nonlocal gravity

History of nonlocal gravity:

- Deser & Woodard, PRL, 2007 (not $\Psi=\Phi$)

$$\Delta\mathcal{L} \equiv \frac{1}{16\pi G} R\sqrt{-g} \times f\left(\frac{1}{\square} R\right),$$

- Maggiore & Mancarella, PRD, 2014 (not dimensionless)

$$S_{\text{NL}} = \frac{1}{16\pi G} \int d^{d+1}x \sqrt{-g} \left[R - \frac{d-1}{4d} m^2 R \frac{1}{\square^2} R \right],$$

- And so on, but none of them satisfy previous 4 requirements

One attempt (Tian, PRD, 2018)

$$\mathcal{L}_G = \frac{1}{2\zeta} \left(R_{\mu\nu} \square^{-1} R^{\mu\nu} - \frac{1}{2} R \square^{-1} R \right),$$

- [zeta]=[kappa]
- -1/2 is for Psi=Phi, based on

$$ds^2 = -c^2(1 + 2\Phi/c^2)dt^2 + (1 - 2\Psi/c^2)d\mathbf{r}^2,$$

$$\square^{-1} R^{\mu\nu} = 0 \text{ in the background}$$

- This assumption has been used in: Koivisto, PRD, 2008; Maggiore & Mancarella, PRD, 2014; Kehagias & Maggiore, JHEP, 2014; Conroy et al, CQG, 2015; etc.

How to set the background?

- Cosmological evolution may give the nonlocal field a nonzero value in the solar system scales.
- Belgacem, Finke, Frassino, and Maggiore, JCAP, 2019 proved this with the McVittie metric.
- And then, Deser2007 and Maggiore2014 (RR model) theories are ruled out by the time-varying G observations. (Belgacem2019)

theory: $|\dot{G}_{\text{eff}}/G_{\text{eff}}| = \mathcal{O}(H_0)$.

observations: $|\dot{G}/G| < 10^{-12} \text{ yr}^{-1} \approx 0.01H_0$

Revisit Tian2018 model

$$\mathcal{L}_G = \frac{1}{2\zeta} \left(R_{\mu\nu} \square^{-1} R^{\mu\nu} - \frac{1}{2} R \square^{-1} R \right),$$

$$ds^2 = -c^2(1 + 2\Phi/c^2)dt^2 + (1 - 2\Psi/c^2)d\mathbf{r}^2,$$

$$U_{00} = f_{10}c^2 + f_{11}(\mathbf{r})c^2, \quad U_{0i} = 0, \quad U_{\mu\nu} \equiv \square^{-1} R_{\mu\nu}$$

$$U_{ij} = \delta_{ij}f_{20} + \delta_{ij}f_{21}(\mathbf{r}),$$

- Results:

Psi=Phi

$$G_{\text{eff}} = \frac{\zeta c^4}{8\pi[-(f_{10} + f_{20})^2 + f_{10} - f_{20} + 1]},$$

$$|\dot{G}_{\text{eff}}/G_{\text{eff}}| = \mathcal{O}(\bar{H}_0).$$

Newtonian approximation in nonlocal Gauss-Bonnet gravity

Capozziello et al, PLB, 20

$$\mathcal{L}_G = \frac{1}{2\kappa} (R - \alpha \mathcal{G} \square^{-1} \mathcal{G}), \quad [\alpha] = \text{length}^4$$

$$\mathcal{G} = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma},$$

$$ds^2 = -c^2(1 + 2\Phi/c^2)dt^2 + a^2(1 - 2\Psi/c^2)d\mathbf{r}^2,$$

$$U = U_0(t) + U_1(\mathbf{r}, t), \quad U \equiv \square^{-1} \mathcal{G}$$

$$H_0/c \approx \partial/(c\partial t) \ll \partial/(a\partial x)$$

$$\gamma \equiv \frac{\Psi}{\Phi} = \frac{c^4/(8\alpha) + 8H^4 + 8H^2\ddot{A} + C_1H/a^3}{c^4/(8\alpha) - 24H^4 + 24H^2\ddot{A} + 16\ddot{A}^2 - 3C_1H/a^3}, \quad \ddot{A} \equiv \ddot{a}/a.$$

$$G_{\text{eff}} = \frac{\kappa c^4}{4\pi(\gamma r_1 + r_2)}, \quad r_1 = \frac{2H^2}{\ddot{A}} + 2 - \frac{48\alpha C_1 H^3}{\ddot{A} a^3 c^4} + \frac{384\alpha H^2 \ddot{A}}{c^4} + \frac{16\alpha C_1 H}{a^3 c^4} - \frac{384\alpha H^6}{\ddot{A} c^4} + \frac{512\alpha H^4}{c^4},$$

$$r_2 = -\frac{2H^2}{\ddot{A}} - \frac{16\alpha C_1 H^3}{\ddot{A} a^3 c^4} - \frac{128\alpha H^6}{\ddot{A} c^4} - \frac{192\alpha H^4}{c^4}.$$

- Generally, $|\dot{G}_{\text{eff}}/G_{\text{eff}}| = \tilde{\mathcal{O}}(H_0)$. $\Psi \neq \Phi$.
- In the de Sitter phase, $\Psi = \Phi$ and G_{eff} is constant.

Summary

- The RT model (Maggiore, PRD, 2014) is the only viable nonlocal gravity theory when confronting observations for now. (Belgacem et al, JCAP, 2019; Tian & Zhu, PRD, 2019)
- Nonlocal gravity may cannot achieve the initial four requirements ($\Psi=\Phi$, dimensionless...), and may cannot solve the coincidence problem.

Thanks